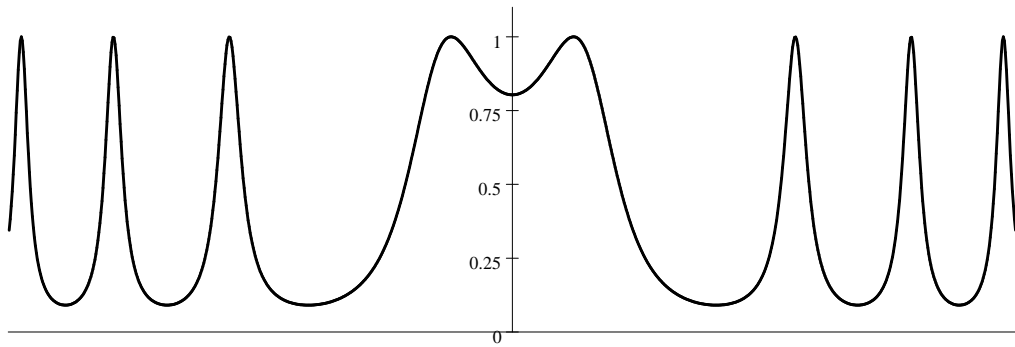


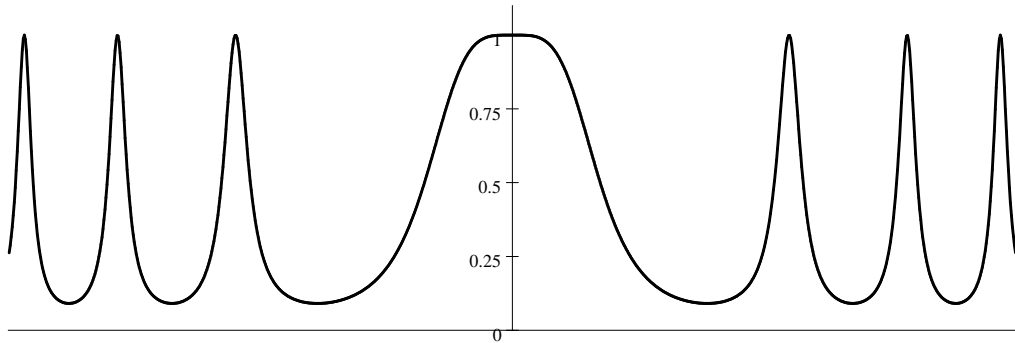
Central wavelength and ring pattern of a Fabry-Perot etalon

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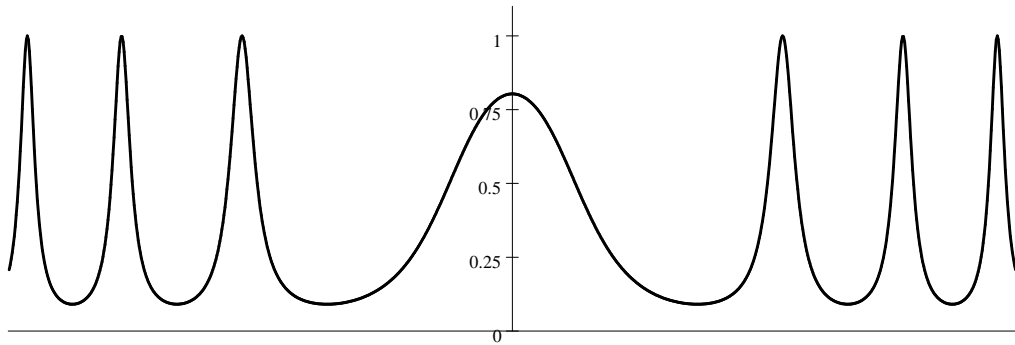
Case 1 : $\lambda_0 > \lambda$, $T = (1 + 10 \sin^2(10.05\pi \cos x))^{-1}$



Case 2 : $\lambda_0 = \lambda$, $T = (1 + 10 \sin^2(10\pi \cos x))^{-1}$



Case 3 : $\lambda_0 < \lambda$, $T = (1 + 10 \sin^2(9.95\pi \cos x))^{-1}$



The central wavelength of a fabry-Perot etalon

The **transmission factor** of the etalon is the ratio, between 0 and 100%, of the transmitted intensity divided by its maximum possible value. This factor T is given by the Airy formula, explained in optics textbooks,

$$T = \frac{1}{1 + F \sin^2 \left(\frac{2\pi ne \cos r}{\lambda} \right)}. \quad (1)$$

The notation is as follows :

- ▶ $F = 4R/(1 - R)^2$, where R is the reflection coefficient of the etalon surfaces
- ▶ n is the refractive index of the gap ($n = 1$ for an air-spaced etalon, $n \simeq 1.6$ for a mica-spaced etalon)
- ▶ e is the gap spacing, expressed with the same unit as the light wavelength λ , thus giving $(2\pi ne \cos r)/\lambda$ in radians.
- ▶ r is the angle of refraction inside the etalon, linked to the angle of incidence i of light on the etalon by Snell's law $\sin r = (\sin i)/n$ (or simply $r \simeq i/n$ if i is small).

A wavelength λ_0 is called a **central wavelength** (CWL) of the etalon if it is one of the wavelengths of maximum transmission under normal incidence, that is $T = 1$ for $\lambda = \lambda_0$ and $i = r = 0$. This means that $\sin^2(2\pi ne/\lambda_0) = 0$ i.e.

$$2ne = k\lambda_0 \quad (2)$$

for some integer k .

When this etalon is lit by a light beam with wavelength λ (e.g. a hydrogen lamp with $\lambda = 0.65628 \mu\text{m}$) its transmission factor is thus

$$T = \frac{1}{1 + F \sin^2 \left(k\pi \frac{\lambda_0}{\lambda} \cos r \right)}. \quad (3)$$

For a given etalon, with fixed values of F , n and e , the accompanying graphs show the variation of T when i (and r) vary. Their maxima give the radii of the **ring pattern** produced by the etalon. For the purpose of illustration we have chosen here $F = 10$, $n = 1$ (air-spaced etalon, so that $i = r$), $k = 10$ and the three examples

$$\text{Case 1 : } \lambda_0 = 1.005\lambda, \text{ Case 2 : } \lambda_0 = \lambda, \text{ Case 3 : } \lambda_0 = 0.995\lambda.$$

The first etalon is tuned to a wavelength slightly $> \lambda$, the second to λ exactly and the third to a wavelength slightly $< \lambda$.

Finding the CWL from the ring pattern

The radii of the ring pattern produced by the etalon correspond to the angles i giving the transmission its maximal value $T = 1$. As the angle i increases from 0, so does the angle r and $k\frac{\lambda_0}{\lambda} \cos r$ decreases, starting from the value $k\frac{\lambda_0}{\lambda}$, close to the

integer k if the CWL λ_0 is close to the lamp wavelength λ . For the m -th ring we have

$$k \frac{\lambda_0}{\lambda} \cos r_m = k - m,$$

where $m \geq 0$ is an integer. If λ_0 is (slightly) greater than λ the first bright ring occurs for $m = 0$, with $\cos r_0 = \lambda/\lambda_0$. This ring is a poorly defined part of the central spot (see the graph for Case 1). If $\lambda_0 \leq \lambda$, the first ring is obtained for $m = 1$ (see Cases 2 and 3).

Thus, leaving aside the case $m = 0$, the ring pattern is defined by

$$\cos r_m = \frac{\lambda}{\lambda_0} \left(1 - \frac{m}{k}\right) \text{ with } m = 1, 2, 3, \dots \quad (4)$$

If the angle of incidence i_m is small enough we may obtain it by the approximate formula $i_m \simeq d_m/2f$ where d_m is the diameter of the m -th bright ring at the prime focus of a camera lens with focal length f . Then $r_m \simeq i_m/n$ by Snell's law, $\cos r_m \simeq 1 - (i_m^2/2n^2)$ and (4) becomes

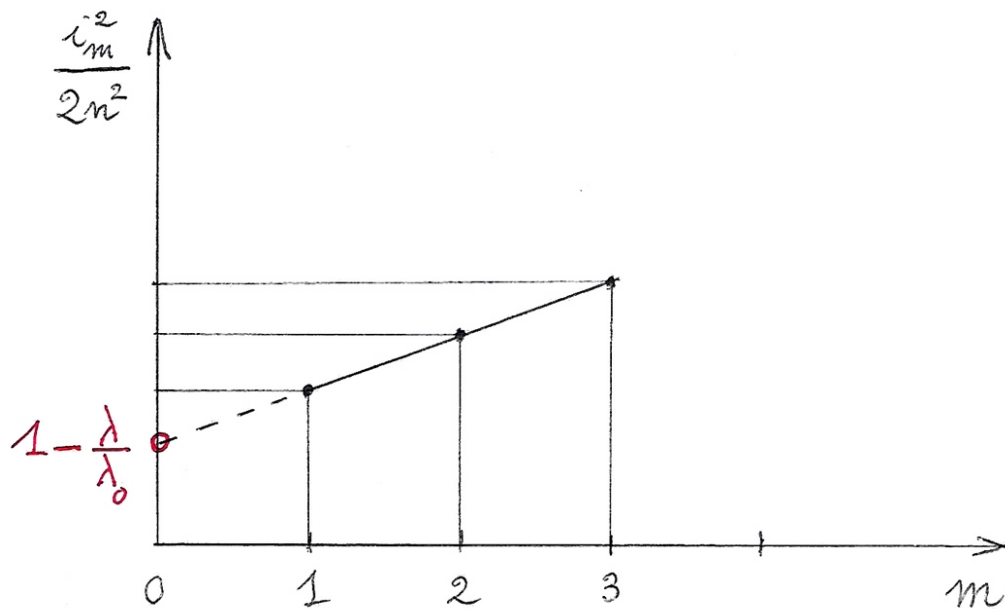
$$\frac{i_m^2}{2n^2} \simeq 1 - \frac{\lambda}{\lambda_0} \left(1 - \frac{m}{k}\right). \quad (5)$$

The graph of $i_m^2/2n^2$ as a function of the ring number $m \geq 1$ is *approximately linear*. When extrapolating it down to $m = 0$, it crosses the vertical axis at $1 - (\lambda/\lambda_0)$ (which may be positive or negative) and this gives (approximately) the etalon CWL λ_0 (see next figure).

Remark. If the above approximations aren't good enough, they may be replaced by exact formulas as follows. Considering an air-spaced etalon ($n = 1$) for the sake of simplicity, we have $i_m = r_m$ and (4) becomes

$$\left(1 + \left(\frac{d_m}{2f}\right)^2\right)^{-1/2} = \frac{\lambda}{\lambda_0} \left(1 - \frac{m}{k}\right) \quad (6)$$

(exact formula) where d_m is the diameter of the m -th ring measured at the focal plane of the camera lens with focal length f . The plot of the left-hand side as a function of $m \geq 1$ is then *exactly linear*. Its extrapolation down to $m = 0$ crosses the vertical axis at λ/λ_0 and this gives the CWL λ_0 . A larger number of rings can now be taken into account for better accuracy.



Approximate measure of the CWL λ_0 of a Fabry-Perot etalon, by extrapolation from the diameter of its first rings.